CS 599 P1: Introduction to Quantum Computation

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PRACTICE WORKSHEET #2

This is the second practice worksheet—it will not be graded and is meant to help you practice the tensor product notation, as well as quantum measurements. I encourage you to work through these problems by **Tuesday, September 23rd**—before the next homework assignment is out. **Solutions will be posted 09/21**.

Background: Tensor product. Let \mathcal{H}_A and \mathcal{H}_B be two Hilbert spaces, e.g. $\mathcal{H}_A = \mathbb{C}^2$ and $\mathcal{H}_B = \mathbb{C}^2$. For any $|\psi\rangle$, $|\psi_1\rangle$, $|\psi_2\rangle \in \mathcal{H}_A$, $|\phi\rangle$, $|\phi_1\rangle$, $|\phi_2\rangle \in \mathcal{H}_B$, and $\alpha, \beta \in \mathbb{C}$, the tensor product \otimes satisfies:

• Linearity in the first argument:

$$(\alpha |\psi_1\rangle_A + \beta |\psi_2\rangle_A) \otimes |\phi\rangle_B = \alpha |\psi_1\rangle_A \otimes |\phi\rangle_B + \beta |\psi_2\rangle_A \otimes |\phi\rangle_B.$$

• Linearity in the second argument:

$$|\psi\rangle_A \otimes (\alpha |\phi_1\rangle_B + \beta |\phi_2\rangle_B) = \alpha |\psi\rangle_A \otimes |\phi_1\rangle_B + \beta |\psi\rangle_A \otimes |\phi_2\rangle_B.$$

• Scalars factor out:

$$(\alpha |\psi\rangle_A) \otimes |\phi\rangle_B = \alpha (|\psi\rangle_A \otimes |\phi\rangle_B).$$

• Zero vector identity:

$$\mathbf{0}_A \otimes |\phi\rangle_B = \mathbf{0}_{AB} = |\psi\rangle_A \otimes \mathbf{0}_B.$$

Problem 1 (Two-qubit tensor products). Recall that a two-qubit state lives in a tensor product space $\mathbb{C}^2 \otimes \mathbb{C}^2$, which has a computational basis consisting of four basis states

$$\{|0\rangle \otimes |0\rangle, |0\rangle \otimes |1\rangle, |1\rangle \otimes |0\rangle, |1\rangle \otimes |1\rangle\}.$$

In this practice exercise, you are asked to express each of the following set of two-qubit states as a complex linear combination of the four basis vectors above.

1. Expand the two-qubit state

$$\left(\frac{1}{\sqrt{2}}\ket{0} - \frac{i}{\sqrt{2}}\ket{1}\right) \otimes \ket{0}.$$

2. Expand the two-qubit state

$$|1\rangle \otimes \left(\frac{\sqrt{3}}{2}|0\rangle + \frac{1}{2}|1\rangle\right).$$

3. Expand the two-qubit state

$$\left(\frac{1}{\sqrt{3}}\left|0\right\rangle+\frac{\sqrt{2}}{\sqrt{3}}\left|1\right\rangle\right)\otimes\left(\frac{1}{\sqrt{2}}\left|0\right\rangle-\frac{i}{\sqrt{2}}\left|1\right\rangle\right).$$

Hint: Try to expand one qubit at a time, using linearity in the tensor product.

Problem 2 (Two-qubit gates). In this exercise, you will practice how to two-qubit gates on the tensor product space $\mathbb{C}^2 \otimes \mathbb{C}^2$. Specifically, you will use the following set of gates:

CNOT
$$|a, b\rangle = |a, a \oplus b\rangle$$
,
CZ $|a, b\rangle = (-1)^{a \cdot b} |a, b\rangle$,
SWAP $|a, b\rangle = |b, a\rangle$.

For brevity, we use denote the basis of the product space $\mathbb{C}^2 \otimes \mathbb{C}^2$ via $\{|0,0\rangle, |0,1\rangle, |1,0\rangle, |1,1\rangle\}$.

1. Apply the CNOT gate to the state

$$\left(\frac{1}{\sqrt{2}}\ket{0} - \frac{i}{\sqrt{2}}\ket{1}\right) \otimes \ket{0}.$$

Expand the resulting state in the basis $\{|0,0\rangle, |0,1\rangle, |1,0\rangle, |1,1\rangle\}$.

2. Apply the Controlled-Z (CZ) gate to the state

$$|+\rangle \otimes |1\rangle$$
, where $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$.

Expand the resulting state in the computational basis.

3. Apply the SWAP gate to the state

$$\frac{1}{\sqrt{2}}(|01\rangle - i|10\rangle).$$

Expand the resulting state.

Problem 3 (Two-qubit measurements) In this exercise, we will consider measurements which take place in either the computational basis $\{|0\rangle, |1\rangle\}$ or the Hadamard basis $\{|+\rangle, |-\rangle\}$, where $|\pm\rangle = \frac{1}{\sqrt{2}}(|0\rangle \pm |1\rangle)$. Consider the following two-qubit entangled state

$$|\Psi^+\rangle = \frac{|01\rangle + |10\rangle}{\sqrt{2}}.$$

- 1. Measure *both* qubits in the computational basis $\{|00\rangle, |01\rangle, |10\rangle, |11\rangle\}$. List each outcome, its probability, and the resulting measurement operators in the form of orthogonal projectors. What is the post-measurement state for each possible outcome?
- 2. Measure only the *first* qubit in the Hadamard basis $\{|+\rangle, |-\rangle\}$, leaving the second qubit unmeasured. For each outcome "+" or "-", list its probability and the measurement operators in the form of orthogonal projectors. What does the post-measurement state look like for each outcome? What is the conditional state of the second qubit?
- 3. Measure the *first* qubit in the computational basis and the *second* qubit in the Hadamard basis simultaneously. For each of the four possible outcomes $\{0, +\}, \{0, -\}, \{1, +\}, \{1, -\}$, list their probability and the corresponding measurement operators in the form of orthogonal projectors. What is the renormalized post-measurement state for each outcome?

Hint: You may find it useful to expand $|\Psi^{+}\rangle$ by replacing the $\{|0\rangle, |1\rangle\}$ basis vectors of the second qubit as linear combinations of the $\{|+\rangle, |-\rangle\}$ basis vectors (see Homework #1).

Recall also that measuring a single qubit of a two-qubit state, in general, results in a re-normalized post-measurement state (see measurement axiom in Lecture 4).

Problem 4 (Entangled states are not product states) Entanglement is a fundamental feature of quantum mechanics: some quantum states of multiple qubits cannot be expressed as a product of single-qubit states. In this exercise, we show that the standard Bell pair is genuinely entangled.

Consider the Bell state

$$|\Phi^+\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle).$$

1. Suppose (for contradiction) that $|\Phi^+\rangle$ can be written as a product of two single-qubit states,

$$|\Phi^{+}\rangle = (\alpha |0\rangle + \beta |1\rangle) \otimes (\gamma |0\rangle + \delta |1\rangle),$$

with complex numbers $\alpha, \beta, \gamma, \delta$ satisfying normalization conditions. Write out the right-hand side explicitly in the computational basis.

- 2. By comparing coefficients of $|00\rangle$, $|01\rangle$, $|10\rangle$, $|11\rangle$, show that no choice of α , β , γ , δ can reproduce the amplitudes of $|\Phi^{+}\rangle$.
- 3. Conclude that $|\Phi^+\rangle$ cannot be expressed as a product state, and therefore is *entangled*.

Problem 5 (Quantum no-cloning theorem) One of the most surprising results in quantum information is that an *arbitrary unknown* quantum state cannot be copied perfectly. This is known as the *no-cloning theorem*. It has profound implications: it protects quantum cryptography, forbids faster-than-light signaling via entanglement, and fundamentally distinguishes quantum information from classical information.

1. Assume that there exists a universal unitary operator U that can clone arbitrary single-qubit states. That is, for any qubit state $|\psi\rangle$ and a fixed blank state $|0\rangle$, the unitary satisfies

$$U(|\psi\rangle\otimes|0\rangle) = |\psi\rangle\otimes|\psi\rangle.$$

2. Check that such a unitary would correctly clone the basis states $|0\rangle$ and $|1\rangle$:

$$U(|0\rangle \otimes |0\rangle) = |0\rangle \otimes |0\rangle, \quad U(|1\rangle \otimes |0\rangle) = |1\rangle \otimes |1\rangle.$$

- 3. Now consider the superposition state $|+\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$. Compute what $U(|+\rangle \otimes |0\rangle)$ would give, assuming *linearity* of U. Compare this to the desired cloned state $|+\rangle \otimes |+\rangle$.
- 4. Show that these two outcomes do not yield the same state. Explain why this means that no such unitary U can exist.